

RECALL: S-L EQ: $\frac{d}{dx} \left(p(x) \frac{d\phi}{dx} \right) + q(x)\phi + \lambda \sigma(x)\phi = 0$, $a < x < b$
 p, q, σ real-valued functions, $p(x) > 0$, $\sigma(x) > 0$
 "weight function"

REGULAR: $\beta_1 \phi(a) + \beta_2 \frac{d\phi}{dx}(a) = 0$, $\beta_3 \phi(b) + \beta_4 \frac{d\phi}{dx}(b) = 0$

LAST TIME:

2. $\phi'' + 4\phi' + 8\phi + \lambda\phi = 0$, $\phi(0) = \phi(L) = 0$

We multiplied by e^{4x} to put this eq. in S-L form.

Solve: char. poly. $r^2 + 4r + (8 + \lambda) = 0$ so $r = -2 \pm \sqrt{-\lambda - 4}$

No nontrivial solutions for $-\lambda - 4 \geq 0$

If $-\lambda - 4 < 0$: $\phi(x) = A e^{-2x} \cos(x\sqrt{\lambda+4}) + B e^{-2x} \sin(x\sqrt{\lambda+4})$

$\lambda + 4 > 0$

Boundary: $\phi(0) = 0 \Rightarrow 0 = A e^0 \cos(0) + 0$
 $0 = A$

$\phi(L) = 0 \Rightarrow 0 = B e^{-2L} \sin(L\sqrt{\lambda+4})$

We need $L\sqrt{\lambda+4} = n\pi$, or $\sqrt{\lambda+4} = \frac{n\pi}{L}$, or $\lambda = \left(\frac{n\pi}{L}\right)^2 - 4$
 for $n = 1, 2, 3, \dots$

Verify S-L Theorems 1-5:

- Eigenvalues are real, and they form an infinite sequence with a smallest, but no largest, eigenvalue
- Eigenfunctions: $\phi_n(x) = e^{-2x} \sin\left(\frac{n\pi}{L}x\right)$, one eigenfunction per eigenvalue
- Completeness: $f \sim \sum_{n=1}^{\infty} a_n \phi_n(x) = \sum_{n=1}^{\infty} a_n e^{-2x} \sin\left(\frac{n\pi}{L}x\right)$

If $f(x) = \sum_{n=1}^{\infty} a_n e^{-2x} \sin\left(\frac{n\pi}{L}x\right)$, then $f(x)e^{2x} = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi}{L}x\right)$

and we can find a_n : $a_n = \int_0^L f(x) e^{2x} \sin\left(\frac{n\pi}{L}x\right) dx$

• Orthogonality: $\int_0^L \phi_n \phi_m \sigma dx = 0$ if $m \neq n$

Here, $\sigma(x) = e^{4x}$: $\int_0^L \left(e^{-2x} \sin\left(\frac{n\pi}{L}x\right) \right) \left(e^{-2x} \sin\left(\frac{m\pi}{L}x\right) \right) e^{4x} dx = 0$ if $m \neq n$

3. Heat equation with non-constant thermal properties.

Suppose $u(x,t) = \phi(x) h(t)$ and separation of variables produces

$$\frac{dh}{dt} = -\lambda h \quad \text{and} \quad \frac{d}{dx} \left[K_0(x) \frac{d\phi}{dx} \right] = -\lambda c(x) \rho(x) \phi$$

(a) $\frac{d}{dx} \left[K_0 \frac{d\phi}{dx} \right] + \lambda \underbrace{c\rho}_{p} \phi = 0$ and $\phi(0) = \phi(L) = 0$ ← Regular S-L equation
 $p(x) = K_0(x)$, $q(x) = 0$, $r(x) = c(x)\rho(x)$

(b) Assume $\phi_n(x)$ are known for $n=1,2,3,\dots$

Then, the Rayleigh Quotient gives:

$$\lambda_n = \frac{-K_0 \phi \frac{d\phi}{dx} \Big|_0^L + \int_0^L \left[K_0 \left(\frac{d\phi}{dx} \right)^2 - 0 \right] dx}{\int_0^L \phi^2 c \rho dx}$$

zero, since $\phi(0) = \phi(L) = 0$

Assume: $K_0(x) > 0$ for all x
 $c(x) > 0$ " "
 $\rho(x) > 0$ " "

$$\lambda_n = \frac{\int_0^L K_0 \left(\frac{d\phi}{dx} \right)^2 dx}{\int_0^L \phi^2 c \rho dx} > 0$$

eigenvalues are positive!

(c) Since $h(t) = ce^{-\lambda t}$,

series solution is $u(x,t) = \sum_{n=1}^{\infty} a_n \phi_n(x) e^{-\lambda_n t}$

exists by S-L theorems

(d) If $u(x,0) = f(x) = \sum_{n=1}^{\infty} a_n \phi_n(x)$

Since eigenfunctions are orthogonal with respect to weight function $c\rho$:

$$\int_0^L f(x) \phi_m(x) c \rho dx = \sum_{n=1}^{\infty} \int_0^L a_n \phi_n(x) \phi_m(x) c \rho dx = 0 \quad \text{unless } m=n$$

$$\int_0^L f(x) \phi_m(x) c \rho dx = \int_0^L a_m \phi_m^2 c \rho dx$$

Solve for a_n :

$$a_n = \frac{\int_0^L f \phi_n c \rho dx}{\int_0^L \phi_n^2 c \rho dx}$$

§5.5

OPERATOR: a function on functions

example: $L(y) = \frac{d}{dx} \left[p \frac{dy}{dx} \right] + qy$

↑
input to L
is a function
 $y(x)$

example: If $y=x^2$:

$$L(x^2) = \frac{d}{dx} \left[p \frac{d}{dx}(x^2) \right] + qx^2$$

$$= \frac{d}{dx} [p \cdot 2x] + qx^2$$

$$= \frac{dp}{dx} \cdot 2x + 2p + qx^2$$

For this operator L , the S-L equation can be written
 $L(\phi) + \lambda \sigma \phi = 0$.

LAGRANGE'S IDENTITY: Let u and v be any functions of x .

Consider: $L(u) = \frac{d}{dx} \left[p \frac{du}{dx} \right] + qu$ and $L(v) = \frac{d}{dx} \left[p \frac{dv}{dx} \right] + qv$

Then: $uL(v) - vL(u) = u \cdot \frac{d}{dx} \left[p \frac{dv}{dx} \right] + quv - v \frac{d}{dx} \left[p \frac{du}{dx} \right] - quv$

Product Rule: $\frac{d}{dx} \left[u \cdot p \frac{dv}{dx} \right] = \frac{du}{dx} \cdot p \frac{dv}{dx} + u \cdot \frac{d}{dx} \left(p \frac{dv}{dx} \right)$

Subtract $-\left(\frac{d}{dx} \left[v \cdot p \frac{du}{dx} \right] = \frac{dv}{dx} \cdot p \frac{du}{dx} + v \frac{d}{dx} \left(p \frac{du}{dx} \right) \right)$

$\frac{d}{dx} \left[u \cdot p \frac{dv}{dx} \right] - \frac{d}{dx} \left[v \cdot p \frac{du}{dx} \right] = 0 + u \cdot \frac{d}{dx} \left(p \frac{dv}{dx} \right) - v \frac{d}{dx} \left(p \frac{du}{dx} \right)$

SAME! ↷

So: $uL(v) - vL(u) = \frac{d}{dx} \left[u \cdot p \frac{dv}{dx} \right] - \frac{d}{dx} \left[v \cdot p \frac{du}{dx} \right]$ **LAGRANGE'S IDENTITY**

Integrate to obtain:

$$\int_a^b (uL(v) - vL(u)) dx = \int_a^b \left(\frac{d}{dx} \left[u \cdot p \frac{dv}{dx} \right] - \frac{d}{dx} \left[v \cdot p \frac{du}{dx} \right] \right) dx$$

$$\int_a^b (uL(v) - vL(u)) dx = \left[u \cdot p \frac{dv}{dx} - v \cdot p \frac{du}{dx} \right]_a^b$$

$\int_a^b (uL(v) - vL(u)) dx = \left[p \left(u \cdot \frac{dv}{dx} - v \frac{du}{dx} \right) \right]_a^b$ **GREEN'S IDENTITY**

$$\int_a^b (uL(v) - vL(u)) dx - \int_a^b (u'v - v'u) dx \Big|_a^b \quad \text{IDENTITY}$$

For most boundary conditions that we encounter (e.g. Regular S-L boundaries), this is zero.

EXAMPLE: Let ϕ_m and ϕ_n be eigenfunctions of a regular S-L problem, corresponding to different eigenvalues. ($m \neq n$)

So: $\phi_n L(\phi_m) + \lambda_m \sigma \phi_m \phi_n = 0$

$-(\phi_m L(\phi_n) + \lambda_n \sigma \phi_n \phi_m = 0)$ subtract

$$\phi_n L(\phi_m) - \phi_m L(\phi_n) + (\lambda_m - \lambda_n) \sigma \phi_n \phi_m = 0$$

Integrate:

$$\int_a^b (\phi_n L(\phi_m) - \phi_m L(\phi_n)) dx = (\lambda_n - \lambda_m) \int_a^b \sigma \phi_n \phi_m dx$$

Regular S-L boundary conditions imply this is zero.

$$0 = (\lambda_n - \lambda_m) \int_a^b \sigma \phi_n \phi_m dx$$

Thus, either $\lambda_n = \lambda_m$, or ϕ_n and ϕ_m are orthogonal!